

Quantum Computing

Mathias Niepert

Homework 1

Exercise 11.1

What is the entropy associated with the toss of a fair coin?

Probabilistic experiment with two possible outcomes 0 and 1
 $\Omega = \{\text{head, tail}\}$ Random variable $X: \Omega \rightarrow \{0, 1\}$. $P(X = 0) = p_0 = 1/2$ and $P(X = 1) = p_1 = 1/2$

$$\begin{aligned} \Rightarrow H(X) &\equiv -\sum_{j=0}^1 p_j \log p_j \\ &= - (1/2 \log(1/2) + 1/2 \log(1/2)) = 1 \end{aligned}$$

What is the entropy associated with the toss of a fair die?

Probabilistic experiment with six possible outcomes $\Omega = \{1, \dots, 6\}$

Random variable X . $P(X = i) = p_i = 1/6 \forall i \in \Omega$

$$\Rightarrow H(X) \equiv -1/6 \sum_{j=1}^6 \log(1/6) = 2.585$$

How would the entropy behave, if the coin or die were unfair?

In both cases the entropy would decrease, because at least one of the possible outcomes would appear with a higher probability than at least one of the others
 \Rightarrow The system contains "less uncertainty" / the entropy of the system is smaller.

Exercise 11.2

The unlikely an event E is, the more information is gained, when it comes out. With an event with probability 1 no information is gained. So it seems appropriate to let the function I a function which takes the probability of the event as input and with: $I(p) \rightarrow 0$ when $p \rightarrow 1$ and $I(p) \rightarrow \pm \infty$ when $p \rightarrow 0$. This is satisfied for $I = k \log(k)$.

A smooth function of probability f must meet the following properties:

1. $\int_D f(x) = 1$ where D is the domain (here: $D = (0, 1] \subset R$).
2. $f(x) \geq 0 \forall x \in D$.

This could be both satisfied for $I = k \log(p)$ and a certain k with $k < 0$.

It is: $k \log(pq) = k (\log(p) + \log(q)) = k \log(p) + k \log(q)$

Exercise 11.6

$$\begin{aligned}(\log(1-p))' &= \frac{-1}{(1-p)\ln(2)} \\(p \log(p))' &= \log(p) + \frac{1}{\ln(2)} \\(p \log(1-p))' &= -\frac{p}{(1-p)\ln(2)} + \log(1-p)\end{aligned}$$

$$\Rightarrow H_{bin}(p)' = (-p \log(p) - (1-p) \log(1-p))' = \log(1-p) - \log(p)$$

$$\begin{aligned}\log(1-p) - \log(p) &= 0 \Rightarrow \\ \log\left(\frac{1-p}{p}\right) &= 0 \Rightarrow p=1/2\end{aligned}$$

$$H_{bin}(p)'' = -\left(\frac{1}{(1-p)\ln(2)} + \frac{1}{p\ln(2)}\right) \text{ which is always } < 0 \text{ for } 0 < p < 1$$

\Rightarrow Only one maximum for $p=1/2$

Show: States of a qubit is 1-1 correspondence with Bloch Sphere

State of a qubit: $|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$

Probability for state $|0\rangle$ when measured: $|\alpha|^2$

Probability for state $|1\rangle$ when measured: $|\beta|^2$

$$\rightarrow |\alpha|^2 + |\beta|^2 = 1$$

Thus we can write:

$$\begin{aligned}|\psi\rangle &= e^{i\varphi_1} \cos\left(\frac{\theta}{2}\right) |0\rangle + e^{i\varphi_2} \sin\left(\frac{\theta}{2}\right) |1\rangle = \\ &e^{i\varphi_1} \left(\cos\left(\frac{\theta}{2}\right) |0\rangle + e^{i\varphi} \sin\left(\frac{\theta}{2}\right) |1\rangle \right); \text{ with } \varphi_2 = \varphi_1 + \varphi\end{aligned}$$

We can ignore the factor of $e^{i\varphi}$ because it has no observable physical effect.

$$\begin{aligned} \Rightarrow |\psi\rangle &= \cos\left(\frac{\theta}{2}\right)|0\rangle + e^{i\varphi}\sin\left(\frac{\theta}{2}\right)|1\rangle = \\ &\cos\left(\frac{\theta}{2}\right)|0\rangle + \cos(\varphi)\sin\left(\frac{\theta}{2}\right)|1\rangle + \sin(\varphi)\sin\left(\frac{\theta}{2}\right) i |1\rangle \end{aligned}$$

This **defines** the (surface of the) Bloch Sphere, as we have three orthonormal base vectors ($|0\rangle, |1\rangle, i|1\rangle$) and, depending on the angles (θ, φ) , the possible coordinates $(\cos(\frac{\theta}{2}), \cos(\varphi)\sin(\frac{\theta}{2}), \sin(\varphi)\sin(\frac{\theta}{2}))$. (By ignoring the factor $e^{i\varphi}$, we "removed" one dimension from our 4-dimensional C^2 space and the $|\alpha|^2 + |\beta|^2 = 1$ condition normed the length of all possible vectors to 1).

$$R^3, \text{ Bloch Sphere: All unit vectors in the z,x plane: } \begin{pmatrix} \sin(\frac{\theta}{2}) \\ 0 \\ \cos(\frac{\theta}{2}) \end{pmatrix}$$

All possible rotations of these vectors round the z-axis:
 $(0 \leq \theta \leq \pi, 0 \leq \varphi \leq 2\pi)$

$$\begin{pmatrix} \cos\varphi & -\sin\varphi & 0 \\ \sin\varphi & \cos\varphi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \sin(\frac{\theta}{2}) \\ 0 \\ \cos(\frac{\theta}{2}) \end{pmatrix} = \begin{pmatrix} \cos\varphi\sin(\frac{\theta}{2}) \\ \sin\varphi\sin(\frac{\theta}{2}) \\ \cos(\frac{\theta}{2}) \end{pmatrix}$$

Every possible value combination of the pair (θ, φ) defines exactly one (of all) possible states of a qubit and exactly one corresponding point on the Bloch Sphere. $(0 \leq \theta \leq \pi, 0 \leq \varphi \leq 2\pi)$

\Rightarrow There is a bijection between the states of a qubit and the points on the Bloch Sphere.

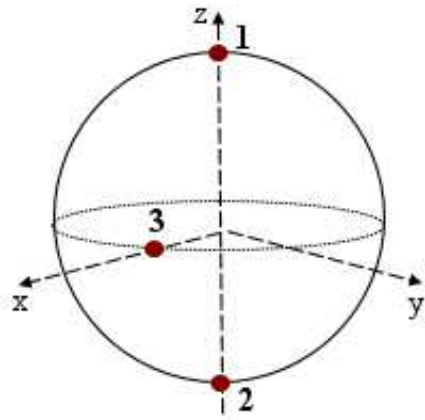


Figure 1: Bloch Sphere

1: $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$

2: $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$

3: $\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$

But there are infinite many states with $|\beta| = \frac{1}{\sqrt{2}}$ and $\alpha = \frac{1}{\sqrt{2}}$